

The background features a dark red to blue gradient with several circular elements. On the left, there is a large circular scale with numerical markings from 140 to 260. Other circles of varying sizes and colors (white, light blue, and dashed) are scattered across the background, some containing arrows or partial arcs.

PRACTICAL APPROACHES AND APPROXIMATIONS II

Laws are like sausages; it is better not to see them being made.

-Otto von Bismarck

Making quality sausage requires watching the process.

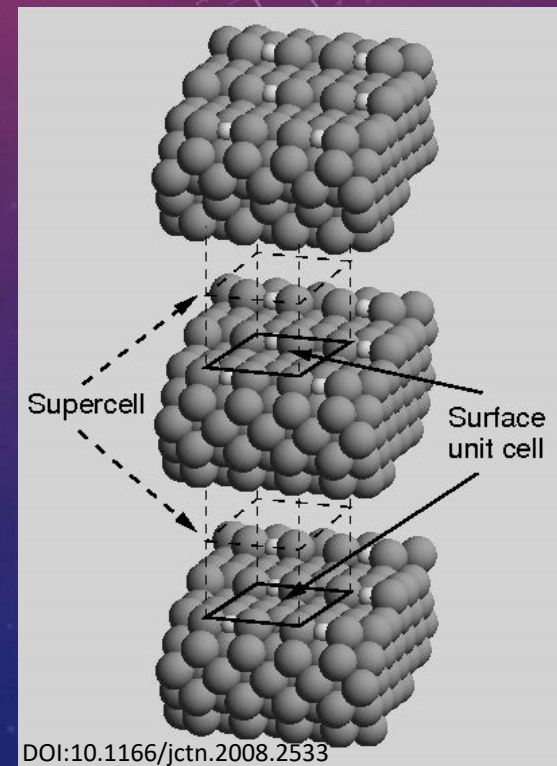
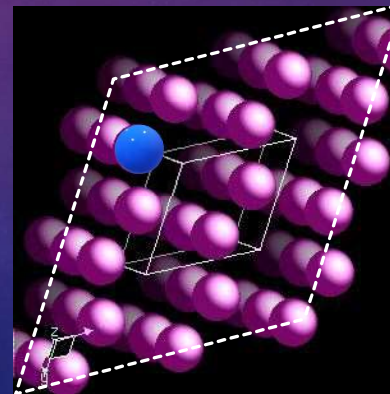
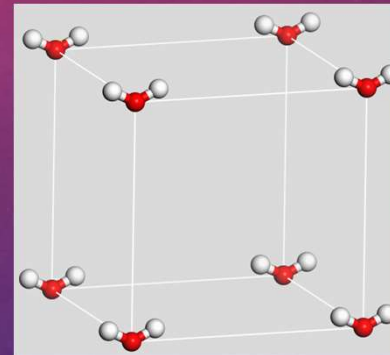
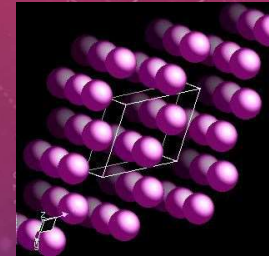
-Tomás Arias, 2023

THIS LECTURE

- Periodicity and aperiodic systems
- Specifying periodic systems
- The cost of periodicity: Bloch's theorem
- Metals and effective electron temperatures

PERIODICITY AND SUPERCELL CONVERGENCE

- Periodicity replaces macroscopic crystal with manageable unit cell
- Aperiodic systems can be embedded in a periodic “supercell”
 - isolated molecules
 - impurities, vacancies, antisite defects
 - surfaces, grain boundaries
 - step edges, dislocations
- Always verify convergence with supercell size (☹)

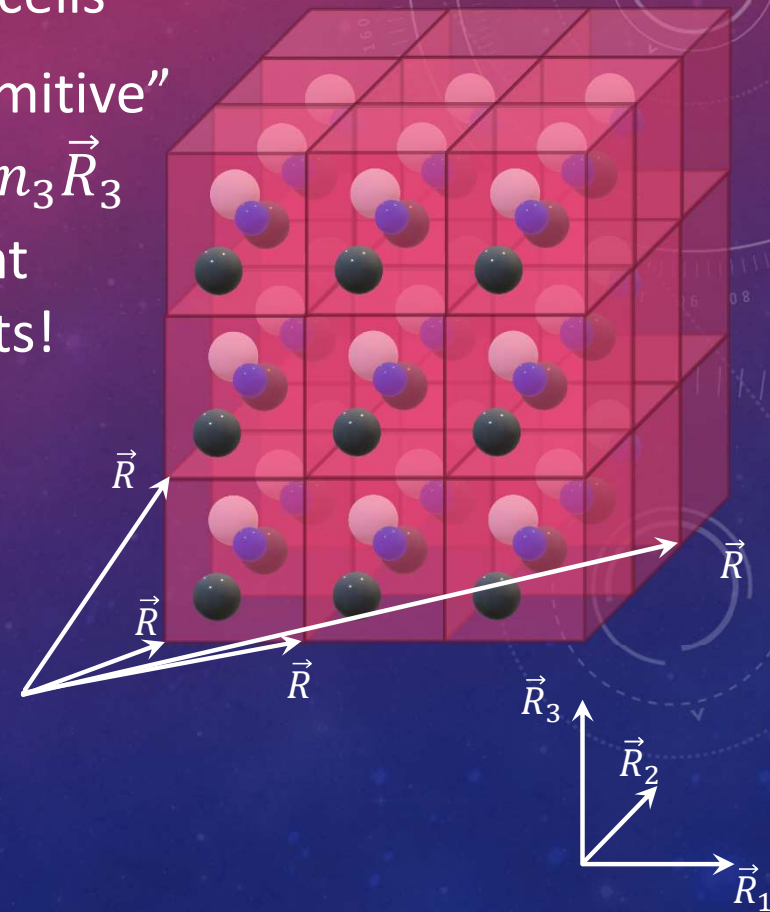
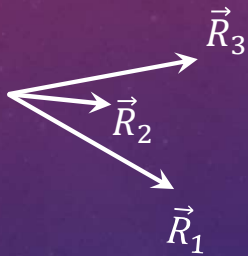


THIS LECTURE

- Periodicity and aperiodic systems
- **Specifying periodic systems**
- The cost of periodicity: Bloch's theorem
- Metals and effective electron temperatures

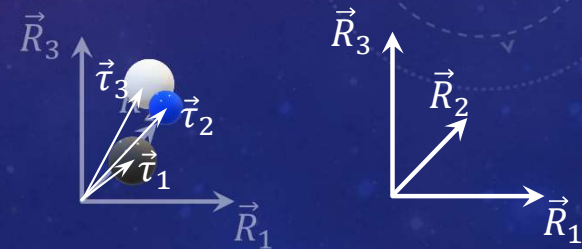
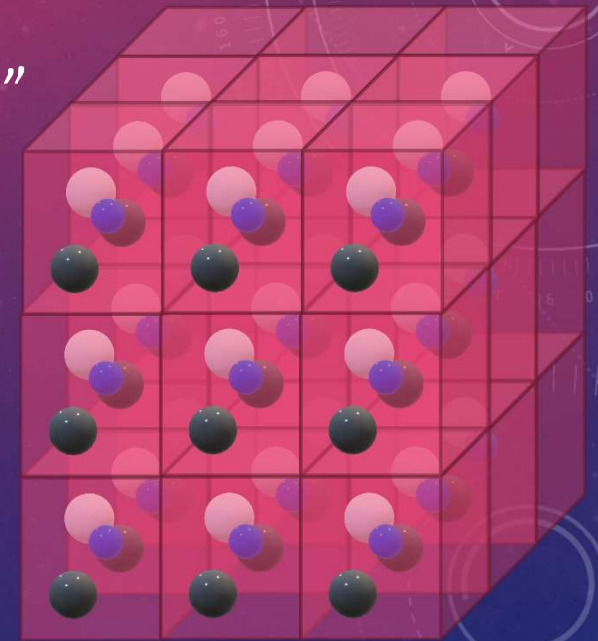
PERIODIC SYSTEMS

- Lattice: Set of locations $\{\vec{R}\}$ of all equivalent unit cells
 - Generally, any \vec{R} can be written in terms of “primitive” lattice vectors $\vec{R}_1, \vec{R}_2, \vec{R}_3$: $\vec{R} = n_1\vec{R}_1 + n_2\vec{R}_2 + n_3\vec{R}_3$
 - Note: $\vec{R}_1, \vec{R}_2, \vec{R}_3$ can be any linearly independent vectors, need not be orthogonal or equal lengths!



PERIODIC SYSTEMS

- Lattice: Set of locations $\{\vec{R}\}$ of all equivalent unit cells
 - Generally, any \vec{R} can be written in terms of “primitive” lattice vectors $\vec{R}_1, \vec{R}_2, \vec{R}_3$: $\vec{R} = n_1\vec{R}_1 + n_2\vec{R}_2 + n_3\vec{R}_3$
- Basis: Locations $\{\vec{\tau}\}$ of atoms within each cell, **relative to the origin of that cell**: $\vec{r}_{\text{atom}} = \vec{R} + \vec{\tau}_{\text{atom}}$
- Wyckoff/“lattice” coordinates: atomic locations expressed as fractions of primitive lattice vectors:
$$\vec{\tau} = x_1\vec{R}_1 + x_2\vec{R}_2 + x_3\vec{R}_3 \text{ (generally } 0 \leq x_\alpha < 1)$$



THIS LECTURE

- Periodicity and aperiodic systems
- Specifying periodic systems
- **The cost of periodicity: Bloch's theorem**
- Metals and effective electron temperatures

BLOCH'S THEOREM I AND K-POINT CONVERGENCE

- Periodicity \Rightarrow all cells are physically identical

$$n(\vec{r} + \vec{R}_\alpha) = n(\vec{r}) \text{ for lattice vectors } \vec{R}_\alpha \ (\alpha = 1, 2, 3)$$

- Wave functions have unobservable phase, so we can only say

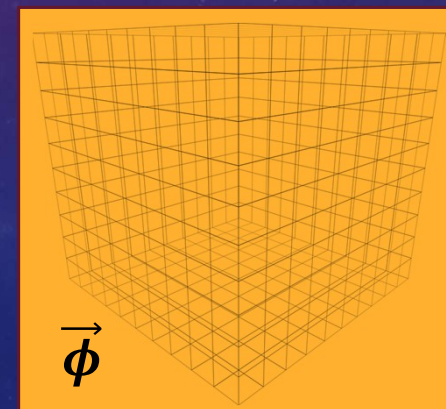
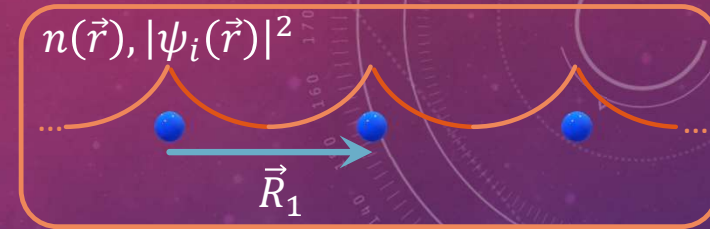
$$|\psi_i(\vec{r} + \vec{R}_\alpha)|^2 = |\psi_i(\vec{r})|^2$$

$$\Rightarrow \psi_i(\vec{r} + \vec{R}_\alpha) = e^{i\phi_\alpha} \psi_i(\vec{r}) \text{ with } \phi_\alpha \in [0, 2\pi)$$

- All $\phi_\alpha \in [0, 2\pi)$ possible and must be included: $\sum_i \leftarrow \iiint_{[0, 2\pi]^3} d^3\phi$

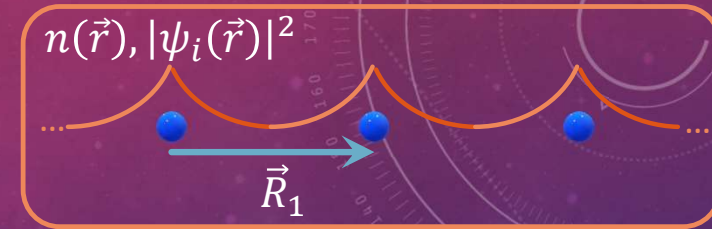
- Integrated by sampling of $[0, 2\pi]^3$ with $N_1 \times N_2 \times N_3$ “k-point” grid

- Always verify convergence with *k-point* sampling (☹)

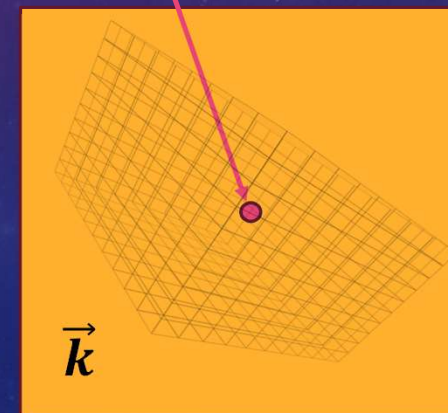


BLOCH'S THEOREM II AND K-POINT SCALING CONVERGENCE

- $\psi_i(\vec{r} + \vec{R}_\alpha) = e^{i\phi_\alpha} \psi_i(\vec{r})$ with $\phi_\alpha \in [0, 2\pi)$
- Consider $\psi_i(\vec{r}) \sim e^{i\vec{k} \cdot \vec{r}}$: $e^{i\vec{k} \cdot (\vec{r} + \vec{R}_\alpha)} = e^{i\vec{k} \cdot \vec{R}_\alpha} e^{i\vec{k} \cdot \vec{r}}$
- Physically, interpret $\phi_\alpha = \vec{k} \cdot \vec{R}_\alpha$ and solve for \vec{k}
- All states $\psi_i(\vec{r})$ have an associated \vec{k}
 \Rightarrow index all with same \vec{k} with $n=1, 2, \dots$: $\psi_i(\vec{r}) \equiv \psi_{n\vec{k}}(\vec{r})$
- **PRO-TIP:** Keeping the same $\Delta\vec{k}_\alpha$ gives same quality
 \Rightarrow upscaling \vec{R}_α upscales $\Delta\phi_\alpha$ and downscales N_α (🤪)
 (bigger supercell, fewer k-points!)



$\psi_{1\vec{k}}(\vec{r}), \psi_{2\vec{k}}(\vec{r}), \psi_{3\vec{k}}(\vec{r}), \dots$



BLOCH'S THEOREM III

- So far, for $\psi_i(\vec{r})$, $i = \{n\vec{k}\}$ and $\phi_\alpha = \vec{k} \cdot \vec{R}_\alpha \dots$

$$\psi_{n\vec{k}}(\vec{r} + \vec{R}_\alpha) = e^{i\vec{k} \cdot \vec{R}_\alpha} \psi_{n\vec{k}}(\vec{r}) \quad [\text{standard form 1}]$$

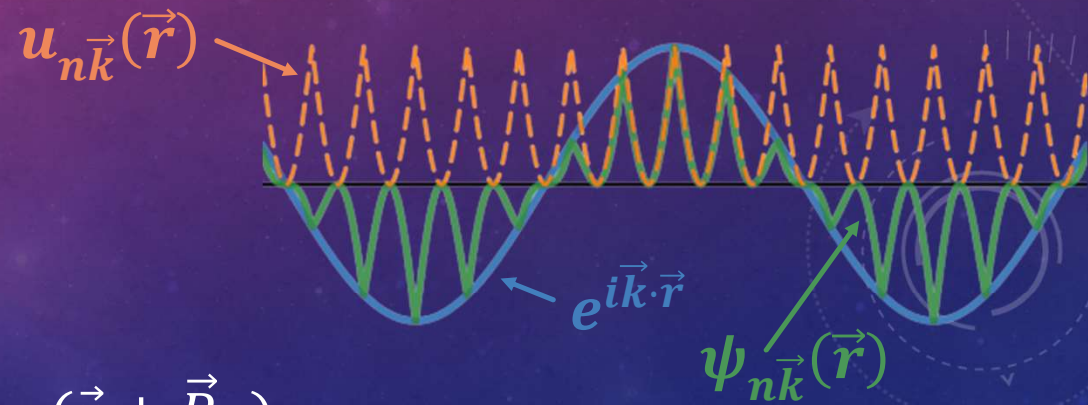
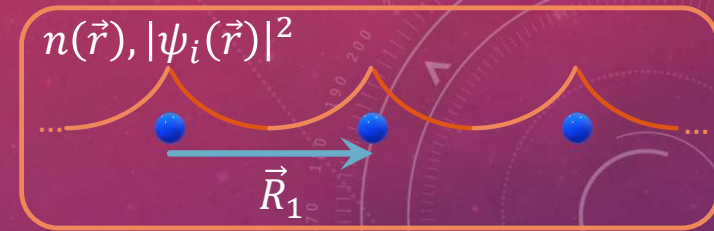
- Alternate form of Bloch's theorem

$$\psi_{n\vec{k}}(\vec{r}) \equiv \underbrace{e^{i\vec{k} \cdot \vec{r}}}_{\substack{\text{wave} \\ \text{envelope} \\ \text{part}}} \underbrace{u_{n\vec{k}}(\vec{r})}_{\substack{\text{periodic} \\ \text{part}}}, \quad u_{n\vec{k}}(\vec{r} + \vec{R}_\alpha) = u_{n\vec{k}}(\vec{r}) \quad [\text{standard form 2}]$$

Proof: $u_{n\vec{k}}(\vec{r}) \equiv e^{-i\vec{k} \cdot \vec{r}} \psi_{n\vec{k}}(\vec{r})$

$$\Rightarrow u_{n\vec{k}}(\vec{r} + \vec{R}_\alpha) = e^{-i\vec{k} \cdot (\vec{r} + \vec{R}_\alpha)} \psi_{n\vec{k}}(\vec{r} + \vec{R}_\alpha)$$

$$= e^{-i\vec{k} \cdot (\vec{r} + \vec{R}_\alpha)} e^{i\vec{k} \cdot \vec{R}_\alpha} \psi_{n\vec{k}}(\vec{r}) = e^{-i\vec{k} \cdot \vec{r}} \psi_{n\vec{k}}(\vec{r}) = u_{n\vec{k}}(\vec{r}) \quad \checkmark$$



THIS LECTURE

- Periodicity and aperiodic systems
- Specifying periodic systems
- The cost of periodicity: Bloch's theorem
- **Metals and effective electron temperatures**

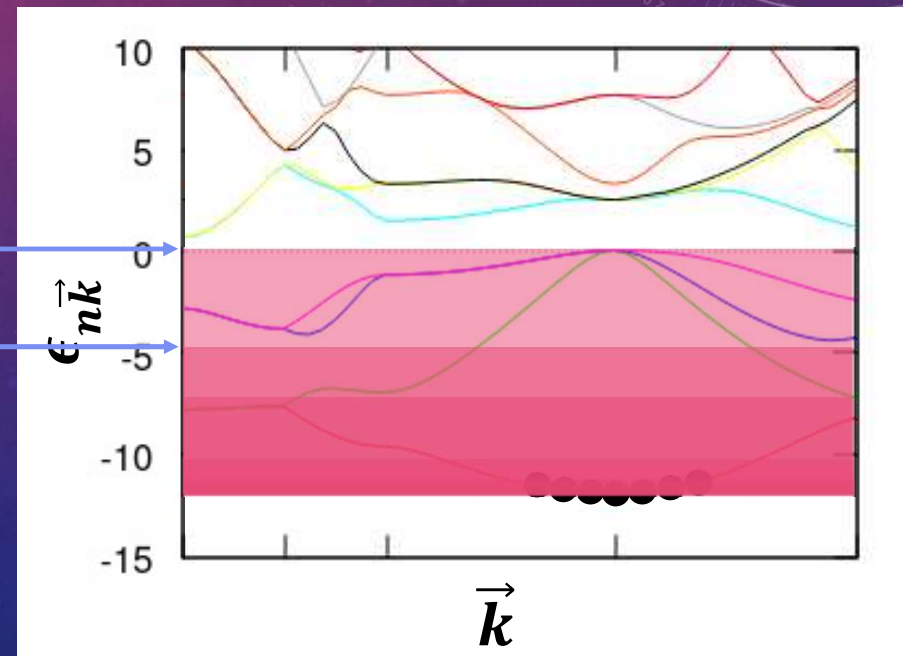
FERMI OCCUPANCY & ELECTRON TEMPERATURE CONVERGENCE

- Electron density (kinetic energy similar story)

$$n(\vec{r}) = \sum_i f_i |\psi_i(\vec{r})|^2 = \sum_{n\vec{k}} f_{n\vec{k}} |\psi_{n\vec{k}}(\vec{r})|^2$$

- At $T=0$, obtain min E_0 by filling lowest energy states first (Aufbau)

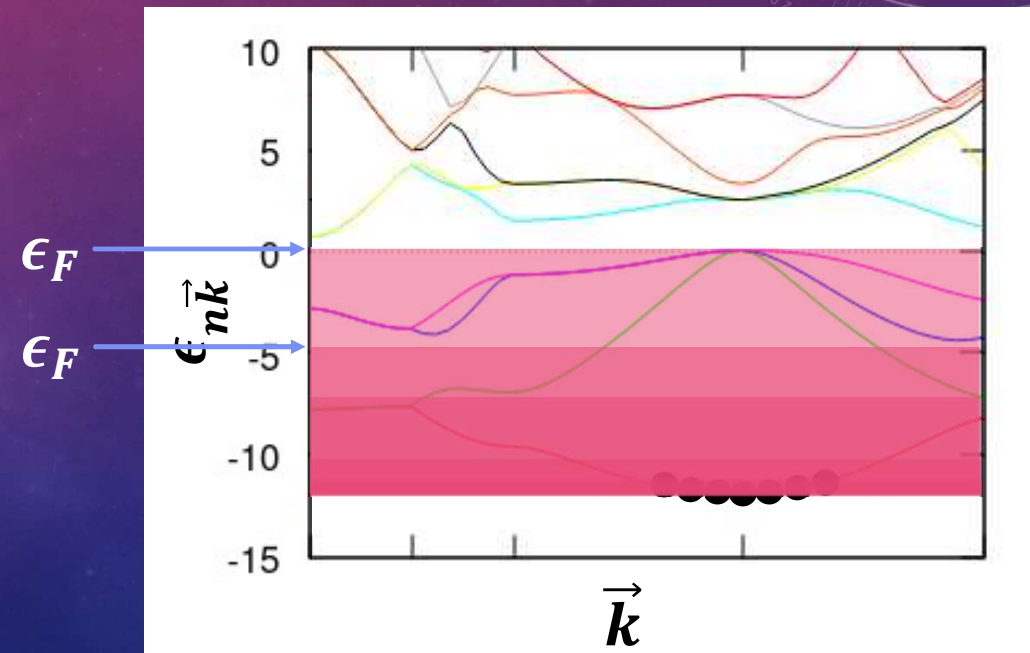
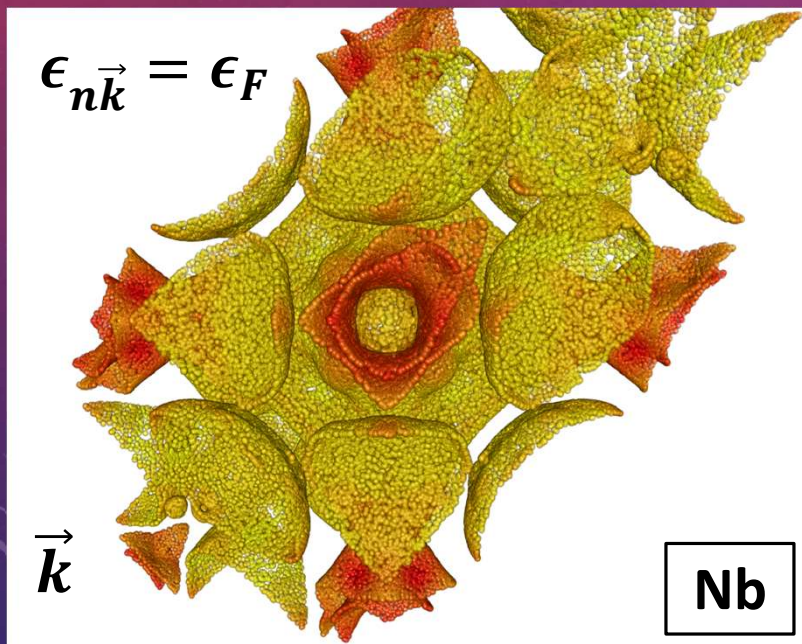
$$f_{n\vec{k}} = \begin{cases} 1, 2 & \text{if } \epsilon_{n\vec{k}} \leq \epsilon_{\text{Fermi}} \\ 0 & \text{if } \epsilon_{n\vec{k}} > \epsilon_{\text{Fermi}} \end{cases}$$



FERMI OCCUPANCY & ELECTRON TEMPERATURE CONVERGENCE

$$f_{n\vec{k}} = \begin{cases} 1, 2 & \text{if } \epsilon_{n\vec{k}} \leq \epsilon_{\text{Fermi}} \\ 0 & \text{if } \epsilon_{n\vec{k}} > \epsilon_{\text{Fermi}} \end{cases}$$

- Semiconductor: completely filled bands: all k=points contribute (🤩)
- Metal: Fermi-level cuts band(s), contributing regions are complicated (😞)

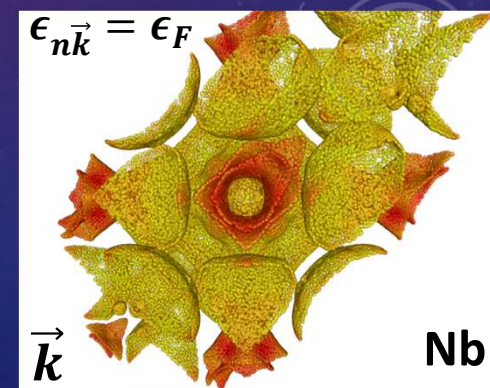
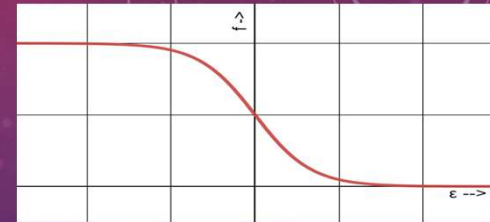


FERMI OCCUPANCY & ELECTRON TEMPERATURE CONVERGENCE

- For metals, contributing regions are complicated (☹️)
- DFT generalizes to $T > 0$!

$$\Rightarrow f_{n\vec{k}} = \frac{1}{e^{(\epsilon_{n\vec{k}} - \mu)/k_B T} + 1} \quad (\text{Fermi-Dirac distribution})$$

- F-D distribution is smooth and easier to integrate (😁)!
 - Higher T needs fewer k-points, with mild impact on results
- \Rightarrow Use elevated T ($\sim 4T_{\text{room}}$), and verify $T \rightarrow 0$ convergence (☹️)



SUMMARY OF ISSUES TO MANAGE

- Choice of approximate functional ($E_{xc}^{LDA}[n(\vec{r})]$, $E_{xc}^{GGA}[n(\vec{r})]$, *etc.*)
 - GGA is usually sufficient (maybe DFT+U for oxides, hybrid for subtle chemistry)
- Iterative convergence (always)
- Plane-wave cutoff E_{cut} (always, once for each new element/pseudopotential)
- Pseudopotential choice (once for each new element, especially exotica)
 - Verify against experimental data for first use
- Supercell (aperiodic systems only)
- k-points (always, but scales inversely with supercell size)
- Electron Temperature (metals only, for every new *compound*)

The background is a gradient from dark purple at the top to dark blue at the bottom. It features numerous out-of-focus circular bokeh lights in shades of purple and blue. On the left side, there are faint, semi-transparent technical diagrams, including a large circular scale with numerical markings from 140 to 260 and several smaller circular elements with arrows and dashed lines.

THANK YOU!

TAA2@CORNELL.EDU